

TABLE II. Observed and calculated values of $\partial T_c/\partial P$.

Element	T_c °K	Atomic volume cm ³	$(C_s - C_n)_{T=T_c}$ mJ deg ⁻¹ mole ⁻¹	$\left(\frac{\partial H_c}{\partial T}\right)_{P=0}^{T=T_c}$ Oe deg ⁻¹	H_0 Oe	$\left(\frac{\partial H_c}{\partial P}\right)_{T=T_c}$ 10 ⁻³ Oe bar ⁻¹	$\left(\frac{\partial T_c}{\partial P}\right)_{H=0}^{\text{calc}}$ 10 ⁻⁶ deg bar ⁻¹	$\left(\frac{\partial T_c}{\partial P}\right)_{H=0}^{\text{obs}}$ 10 ⁻⁶ deg bar ⁻¹
V	5.03 ^a	8.34	69.4 ^a	-455	1310 ^a	4.1±0.3 ^b 2.0±0.2 ^c	9.0 4.4	11±3
Nb	9.17 ^d	10.80	140 ^d	-421±4	1944 ^d	-1.2±0.3 ^e	-2.85	0±3
Ta	4.39 ^e	10.83	42.2 ^e	-334±2	825 ^f	-0.8±0.3 ^e	-2.4	-2.6±1.0 ^f

^a Reference 5.^b Reference 10.^c Reference 8.^d Reference 7.^e Reference 6.^f Reference 9.

in Table II. In order to calculate values of $\partial T_c/\partial P$ using the Maxwell thermodynamic relationship,⁴

$$\left(\frac{\partial T_c}{\partial P}\right)_{H=0} = -\left(\frac{\partial H_c}{\partial P}\right)_{T=T_c} \left(\frac{\partial H_c}{\partial T}\right)_{P=0}^{-1}, \quad (3)$$

we express the measured values⁵⁻⁷ of $C_s - C_n$ in terms of $(\partial H_c/\partial T)_{T=T_c}$ using the Rutgers relationship,⁴

$$(C_s - C_n)_{T=T_c} = \frac{VT_c}{4\pi} \left(\frac{\partial H_c}{\partial T}\right)_{P=0}^2. \quad (4)$$

The values of $\partial H_c/\partial T$, given in Table II, derived in this manner are in good agreement with values obtained from directly measured critical-field curves for vanadium⁵ and tantalum,⁶ but not for niobium.⁷

Using the thermodynamic relationship (4) we have calculated values of $(\partial T_c/\partial P)_{H=0}$, and these are compared in Table II with our observed values. Table II also includes the results for tantalum; $(\partial T_c/\partial P)_{H=0}$ was determined for this element by Hinrichs and Swenson.⁹ The sign of $(\partial T_c/\partial P)_{H=0}$ obtained for vanadium agrees with that predicted from the thermal-expansion data. The observed magnitude is in better agreement with the value calculated from the thermal-expansion data of Müller and Rohrer,¹⁰ rather than the value determined from the data of White.⁸ The calculated value of $(\partial T_c/\partial P)_{H=0}$ for niobium is about the limit of our experimental sensitivity and is, therefore, not inconsistent with the zero pressure dependence observed. The experimental results of Hinrichs and Swenson⁹ are also in good agreement with the calculated value.

The effect of applying pressure to a superconductor, until recently, had always been associated with an observed decrease in the superconducting transition temperature.¹¹ However, a number of superconductors

(Zr,¹² La,¹³ U,¹⁴ and V¹⁵) have now been found to exhibit a positive $\partial T_c/\partial P$. We may attempt to understand this difference in sign of the pressure dependence of the superconducting transition temperature by considering the volume derivative of the BCS¹⁶ relationship,

$$T_c = 0.85\Theta_D \exp(-1/A), \quad (5)$$

with $A = N(0)V$, where $N(0)$ is the density of electron states at the Fermi surface and V is the attractive electron-electron interaction parameter. Differentiation of (5) with respect to volume gives

$$\frac{\partial \ln T_c}{\partial \ln v} = \varphi \ln\left(\frac{0.85\Theta_D}{T_c}\right) - \gamma_G, \quad (6)$$

where $\varphi = \partial \ln A/\partial \ln v$ and γ_G , the Grüneisen constant, represents the volume dependence of the phonon spectrum. Rewriting $\partial \ln T_c/\partial \ln v$ in terms of $\partial T_c/\partial P$ we have

$$\frac{\partial T_c}{\partial P} = -|K|T_c \left\{ \varphi \ln\left(\frac{0.85\Theta_D}{T_c}\right) - \gamma_G \right\}, \quad (7)$$

where K is the compressibility.

The pressure dependence of the phonon spectrum is such as to increase T_c and will be roughly the same for all elements since γ_G has, in general, values between 1 and 3. Since $\ln(0.85\Theta_D/T_c)$ lies in the range 2.5 to 6.5 for most superconductors the sign and magnitude of $\partial T_c/\partial P$ is determined by φ . Rohrer¹⁷ has pointed out that for nontransition metal superconductors φ is roughly constant and equal to 2.5 ± 0.5 . However, when we consider the behavior of the transition metal superconductors there is considerable variation both in the magnitude and the sign of φ .^{18,19} Olsen and his co-

¹² N. B. Brandt and N. I. Ginzburg, Zh. Eksperim. i Teor. Fiz. 46, 1212 (1964) [English transl.: Soviet Phys.—JETP 19, 823 (1964)].

¹³ W. E. Gardner and T. F. Smith, Phys. Rev. 138, A484 (1965).

¹⁴ T. F. Smith and W. E. Gardner, Phys. Rev. 140, A1620 (1965).

¹⁵ Present work.

¹⁶ J. Bardeen, L. N. Cooper, and J. R. Schrieffer, Phys. Rev. 108, 1175 (1957).

¹⁷ H. Rohrer, Helv. Phys. Acta 33, 675 (1960).

¹⁸ J. L. Olsen, E. Bucher, M. Levy, J. Müller, E. Corenzwit, and T. Geballe, Rev. Mod. Phys. 36, 168 (1964).

¹⁹ E. Bucher, J. Müller, J. L. Olsen, and G. Palmy, Phys. Letters 15, 303 (1965).

⁹ C. H. Hinrichs and C. A. Swenson, Phys. Rev. 123, 1106 (1961).

¹⁰ J. Müller and H. Rohrer, Helv. Phys. Acta 31, 289 (1958).

¹¹ An exception to this generalization is thallium which shows a slight increase in T_c for applied pressures up to 2 kbar [Ref. 1; J. Hatton, Phys. Rev. 103, 1167 (1956); and I. D. Jennings and C. A. Swenson, Phys. Rev. 112, 31 (1958)]. Further application of pressure then causes T_c to decrease. Jennings and Swenson have explained this behavior as a consequence of the highly anisotropic nature of the physical properties of thallium.